Structure of the upper mantle in the north-western and central United States from USArray S-receiver functions

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Abstract. We used more than 40 000 S-receiver functions recorded by the USArray project to study the structure of the upper mantle between the Moho and the 410 km discontinuity from the Phanerozoic western United States to the cratonic central US. In the western United States we observed the lithosphere–asthenosphere boundary (LAB), and in the cratonic United States we observed both the mid-lithospheric discontinuity (MLD) and the LAB of the craton. In the northern and southern United States the western LAB almost reaches the mid-continental rift system. In between these two regions the cratonic MLD is surprisingly plunging towards the west from the Rocky Mountain Front to about 200 km depth near the Sevier thrust belt. We interpret these complex structures of the seismic discontinuities in the mantle lithosphere as an indication of interfingering of the colliding Farallon and Laurentia plates. Unfiltered S-receiver function data reveal that the LAB and MLD are not single discontinuities but consist of many small-scale laminated discontinuities, which only appear as single discontinuities after longer period filtering. We also observe the Lehmann discontinuity below the LAB and a velocity reduction about 30 km above the 410 km discontinuity.

1 Introduction

Lithospheric plates, including thick old cratons, translate over thousands of kilometers over the viscous mantle. However, relatively little is still known about the internal structure of cratons and the transition between the craton and the convecting mantle. Even after more than half a century since the general acceptance of plate tectonic theory, the lithosphere–asthenosphere boundary (LAB) below cratons is still thought to be “elusive” (Eaton et al., 2009) and an additional velocity drop frequently observed in seismic data in the shallow cratonic lithosphere (mid-lithospheric discontinuity, or MLD) is referred to as “enigmatic” (Karato, 2012). These descriptions not only apply to the petrophysical properties of the LAB and MLD but are also a result of still inadequate seismological imaging. The lithosphere–asthenosphere system was originally a mechanical definition (Barrell, 1914) and is not a seismic definition. However, we are using the name LAB here for seismic velocity reductions observed near 200 km depth in cratons and near 100 km depth in oceans and Phanerozoic regions by tomography and receiver functions (e.g., Yuan and Romanowicz, 2010). Tomography is not directly sensitive to discontinuities, and therefore the transition from lithosphere to asthenosphere is derived from the velocity–depth functions or its vertical gradients (e.g., Yoshizawa, 2014). Only a few tomography studies observed in cratons a shallow low-velocity zone near 100 km depth, which could be related to the MLD (e.g., Lekic and Romanowicz, 2011).

The alternative view of the deep structure of cratons, independent of the lithosphere–asthenosphere model, defines a tectosphere with the keel of the cratons reaching down to about 400 km (Jordan, 1975). This model is based on different versions of the tomography method (see Jordan and Paulson, 2013, for a summary). The tectosphere may be decoupled from the convecting mantle by a low-viscosity layer directly above the 410 km discontinuity. A possibly related velocity reduction above the 410 is observed globally by Tauzin et al. (2010).

The LAB and MLD are mainly observed with receiver functions. Summaries of their global distribution are given...
by Rychert and Shearer (2009), Rychert et al. (2010), Fischer et al. (2010) and Kind et al. (2012). Since the MLD and in some places the cratonic LAB are relatively well observed with seismic body waves, both discontinuities must be relatively sharp. Rychert et al. (2007) deduced from P-receiver functions a sharpness of 11 km for the LAB at the east coast of the US. Li et al. (2007) obtained a sharpness of about 20 km in the western United States using S-receiver functions. Rychert et al. (2007) concluded that the sharpness of the LAB excluded temperature variation as the single cause of the LAB. Karato (2012) suggested the grain-boundary sliding model as the cause of the MLD and LAB. This model predicts a strong MLD and weak LAB below cratons. According to Yuan and Romanowicz (2010) the North American craton consists of two layers with different chemistry and anisotropy, with the upper layer reaching to about 100 km depth being the Archean lithosphere. Selway et al. (2015) reviewed the mechanisms which could cause the MLD.

It would be difficult to show the existence of the MLD in Phanerozoic regions since here the LAB would also be expected at a similar depth. Therefore it is interesting to study the structure of the MLD and LAB at continental collision zones as in the western US. Abt et al. (2010) and Kumar et al. (2012a, b) observed in almost the entire United States only a shallow negative discontinuity near 100 km depth. This discontinuity varies to some extent in depth but nowhere does it reach 200 km depth. In the western United States and at the east coast this signal was called the LAB while in the central cratonic United States it was called the MLD by Abt et al. (2010), and the LAB by Rychert and Shearer (2009) and Kumar et al. (2012a, b). Levander and Miller (2012) have mapped the Phanerozoic LAB in the western US. More detailed regional studies in the western United States are published by Rychert et al. (2005, 2007), Li et al. (2007), Hopper et al. (2014), Hansen et al. (2013), Foster et al. (2014) and Lekic and Fischer (2014). Rare observations of the cratonic LAB near about 200 km depth have been obtained by Foster et al. (2014) in the US. Similar observations in other cratons have been obtained in Canada (Miller and Eaton, 2010), Scandinavia (Kind et al., 2013) and South Africa (Sodoudi et al., 2013).

The structure of the mantle lithosphere in western North America was formed by the collision of the Farallon plate with the Laurentia craton and was first resolved by to-mographic studies. The subducted Farallon plate is visible in the upper and lower mantle, even below the eastern United States (Grand, 1994; Schmandt and Lin, 2014). The collision with the Precambrian North American craton about 50 million years ago during the Laramide orogeny tore and broke the Farallon plate. An example is the “big break” in the western United States of Sigloch et al. (2008) and Sigloch (2011) and references therein, or the vertical high-velocity “curtain” near the longitude of Yellowstone (Schmandt and Humphreys, 2011). The boundary of the craton follows approximately the Rocky Mountain Front (see for example Yuan et al., 2011, and Zheng and Romanowicz, 2012).

2 Methodology

We use the S-receiver function technique (meaning S-to-P conversions) to image seismic discontinuities between the crust–mantle boundary (Moho) and the seismic discontinuity at 410 km depth (see, e.g., Yuan et al., 2006, or Kind et al., 2012, for a description of the technique). The receiver function method determines the response of the Earth structure below a seismic station. Teleseismic waves arriving at a station are scattered by the underlying discontinuities, causing conversions and multiple reflections, which lead to images of the layered structure below the station.

The most important step in receiver function processing is stacking of many seismic traces in order to enhance the weakly converted waves. The simplest approach is to align many records for a given station with respect to a main phase, for example S, after amplitude and sign normalization, and sum these traces. The summation is performed for each component separately (vertical, radial and transverse). We rotated the components by theoretical back azimuth and incidence angle of the S phase and obtained approximately the P, SV and SH response of the medium below the station. As the travel-time moveout between the main phase and the scattered phase depends on epicentral distance, this kind of summation can, without a moveout correction, only be done in narrow epicentral distance windows (Shearer, 1991). A distance moveout correction permits summation over larger distance ranges (Yuan et al., 1997). However, a velocity model is required. Since the moveout correction can only be made for one type of phase at a time (for example P-to-S conversions, or multiples), signals traveling with different slownesses will be canceled by stacking.

Traditionally, deconvolution is used to equalize (and approximately remove) the source signals of the different earthquakes before stacking. For example, in P-receiver functions, a window around the P signal on the vertical component is used to deconvolve the radial component. We should also mention that, theoretically, deconvolution is not required in the receiver function technique. Summation of plain seismo-grams leads to similar results (see Kumar et al., 2010; Bodin et al., 2014). Here we applied deconvolution since we obtained better signal-to-noise ratios. Another important step in receiver function processing is the migration from the time domain into the depth domain using a known velocity model. For migration, the seismic energy is back-projected along the ray path within a given model and stacked, assuming that the energy is distributed in the Fresnel zone (Jones and Phinney, 1998; Kosarev et al., 1999; Yuan et al., 2006). The one-dimensional IASP91 reference model is used for moveout correction and migration. Both moveout correction and migration are relatively insensitive to the model used.
We use S-receiver functions in the present study, which have a significant advantage over P-receiver functions for upper mantle studies. In S-receiver functions, the direct conversions arrive before the S signal while the crustal multiples arrive after the S signal. In P-receiver functions, both direct conversions and multiples arrive after the P signal. Multiples in P-receiver functions frequently overwhelm direct conversions and make it difficult to identify the true structure. However, S-receiver functions have other problems which need to be considered. In the next section we include a discussion of some problems of the interpretation of the wave field of S precursors using the example of USArray data.

3 Data

We obtained the data from the open-access IRIS archive in Seattle, Washington (www.iris.edu). Most data are provided by the USArray project (www.usarray.org), which is a continent-wide temporary mobile network with a spacing of about 70 km between stations. Stations recorded on average for about 2 years at one site before they were moved to another site. The locations of the seismic stations used in this study are shown in Fig. 1. The distribution of the epicenters of the earthquakes used is shown in Fig. 2.

We have inspected manually over 200,000 records from events with magnitudes greater than 5.7 within the epicentral distance range of 60–85°. Conditions for the selection of traces were a signal-to-noise ratio of at least two in the original broadband S signal on the SV component, a good approximation of a spike by the deconvolved SV signal, low energy at the time of the spike on the P component and no obviously disturbing signals before the S arrival on the P component. This procedure appears to be robust, since several persons participated in selecting data along these lines without a visible personal influence.

Finally more than 40,000 records have been selected for our study. Next we will discuss the wave field of the S pre-
Figure 3. (a) Display of binned S-receiver functions as a function of the epicentral distance. Each bin contains more than one thousand traces. Precursors of the S phase from S-to-P conversions in the upper mantle are marked with dashed black lines (410 conversion at the 410, Lehmann–Lehmann discontinuity, LVZ1, LVZ2 conversions at velocity reductions between Moho and Lehmann, Moho conversion at the Moho). LVZ2 and Lehmann are more clear in Fig. 4; they are marked here for comparison. Additional theoretical travel-time curves of ScS, SKS, S-to-P conversions of ScS and SKS at the 410 and 660 discontinuities and at a possible discontinuity at 580 km depth (ScS410p, ScS580p, ScS660p, SKS410p, SKS580p, SKS660p), crustal multiples (SPmp) and SPn are marked in green. (b) Ray paths of Sp, SPmp and SPn.

Figure 4. Display of binned S-receiver function traces along a (a) west–east and (b) south–north line. LVZ1 is the LAB in the western United States and the MLD in the central US. Details of LVZ1 and LVZ2 are shown in the narrower profiles in Figs. 8 and 9. “Lehmann” indicates the bottom of the asthenosphere, and LVZ410 marks a velocity drop above the 410 km discontinuity. Note that this discontinuity is only weakly observed at the eastern end of the line. “410” is the discontinuity at 410 km depth. Although these discontinuities are greatly averaged in this display, they appear very clearly.

The back azimuth of the sources is disregarded. The boxes are aligned along west–east and south–north profiles. West–east profiles are shown in Fig. 4a. The box size is 0.5° longitude and extends in the south–north direction over the entire array. South–north profiles are shown in Fig. 4b. The box size is 1° latitude and extends in west–east direction also over the entire array. In Fig. 5 the traces are shown from the cratonic part of the network (east of 110° W longitude) as a function of the back azimuth. There is no overlapping of windows. Neighboring stacked traces therefore do not contain any common traces.

There are obviously several seismic phases visible in Figs. 3–5. Seismic phases marked red (positive) are caused by a discontinuity with downward increasing velocity. Blue phases (negative) mark downward decreasing velocity. All phases converted from the direct S phase to P are marked with black labels in Figs. 3–5. These are the conversions from the Moho, two blue phases labeled LVZ1 and LVZ2, the conversion from the discontinuity at 410 km depth (marked...
Figure 5. Display of binned S-receiver functions as a function of the back azimuth (see Fig. 2) of the source of each record, independent of epicentral distance. Only stations on the cratonic part of the United States (east of 110°W) are used. The same phases as in Fig. 4 are marked. Most phases do not show any clear dependence on the back azimuth. The only exception might be the 410 discontinuity, which is strongest for sources in the northwest quadrant.

"410"), a red discontinuity marked “Lehmann" and another blue discontinuity following closely the 410 signal and marked “LVZ410". The Moho is not the focus of our study. The LVZ410 is observed in P- and S-receiver functions by a number of authors in different parts of the world (e.g., by Schaeffer and Bostock, 2010, in northwestern Canada; by Vinnik et al., 2010, in California and globally by Tauzin et al., 2010). It is interpreted by the presence of water causing partial melt. Jordan and Paulson (2013) discuss the role of this discontinuity for decoupling the thick continental tectosphere (which extends below the LAB) from the convecting mantle. The Lehmann discontinuity is widely considered as the bottom of the asthenosphere. A global study of the Lehmann discontinuity is given by Gu et al. (2001).

Theoretical travel-time curves of S precursors at the above mentioned discontinuities are marked by dashed black lines in Fig. 3a. They are computed using the IASP91 global reference model with three negative discontinuities (LVZ1, LVZ2 and LVZ410 at 90, 170 and 380 km depth, respectively) and one positive discontinuity (Lehmann at 270 km) added. The Moho is set at 35 km depth. LVZ1 and the 410 are clearly observed in Figs. 3–5. The LVZ2 and the Lehmann discontinuity are better observed in Fig. 4.

Additional seismic phases, besides the ones discussed so far, are only visible in Fig. 3a but not in Figs. 4 and 5. Theoretical travel-time curves have been computed to explain these phases. They are marked in green. We see clear negative SKS660p and ScS660p phases from the 660 km discontinuity cutting through all other phases prior to the S arrival. These phases are strong at 68–73° and 75–79° epicentral distances and 10–30 s precursor time, and at 75–80° epicentral distance and 40–50 s precursor time, where they cut through the LVZ2 and the LVZ410 signals. However, the signals are caused by the SKS and ScS conversions at the 410 and 660 discontinuities and have very different slownesses than the S phase. This is the reason why they are canceled out in the moveout-corrected and stacked signals in Figs. 4 and 5. There are also surprisingly clear phases after the S signal in Fig. 3a. They are the crustal multiples SPmp and SPn below the stations (see the ray path of SPn in Fig. 3b), which are so far not much used to infer information about the P velocity below the stations.

There have been concerns that higher-order P multiples could influence S precursors on the P component (Bock, 1994; Wilson et al., 2006). However, we do not see such phases in the complete observed precursor wave field in Fig. 3a. The reason is probably that the upper mantle of the Earth is too heterogeneous to permit the efficient propagation of higher-order P multiples. In Fig. 6 we have computed S precursors for the one-dimensional IASP91 Earth model (vertical component). We used the version of the reflectivity method by Kind (1985), which is an extension of the original method by Fuchs and Muller (1971) for different source and receiver structures. This version avoids P multiples in the the-
Figure 7. Theoretical seismograms for different structures at the source and receiver sides. (a) Moho is at 50 km depth at the source side, (b) no crust at the source side. Moho depth is 40 km at the receiver side in both cases. The source crust causes a negative precursor before the receiver side S-to-P conversion at the Moho (see PmSmp ray path in c), which could be mistaken as an indication of a velocity drop in the mantle below the Moho. Such phases will only be a problem in interpretations of single data traces. In receiver function processing of real data this should not be a problem since source side models and source depths are different for each record and these effects are erased by summation. (c) Ray paths of Sm and PmSm.

oretical seismograms, which cause a high noise level in front of the S phase making it difficult to obtain good S precursors. An approximated spike was used as the source time function. The strong SPn and Sppm phases after S agree well with the data in Fig. 3a. A disagreement between computed and observed seismograms is the source-side Pn in front of the S signal (Fig. 6), which is not in the data (Fig. 3a). This is a phase traveling as P along the surface on the source side and continuously radiating S waves downward. At the receiver side these phases are converted back to P and travel again horizontally along the surface. These phases are not observed in the stacked S-receiver functions (Fig. 3a) because of stacking of records from many regions with different structures and source depths. The same phases marked green in the data (Fig. 3a) are also marked green in the theoretical seismograms in Fig. 6. Note that the IASP91 model has neither an upper-mantle low-velocity zone, the Lehmann discontinuity, nor the negative discontinuity above the 410. Therefore these phases are not computed.

In order to point out another possible source of disturbing S precursors, we have computed theoretical seismograms similar to Fig. 6 for different crustal models at the source and receiver sides (see Fig. 7). A narrower slowness integration window was used, which excludes source-side Pn. There is a 50 km thick crust at the source side in Fig. 7a and no crust is included at the source side in Fig. 7b. Receiver-side crust is in both cases 40 km thick. The rotated L component is shown, which carries practically only P energy. We see ScS and SKS conversions at the 410 and 660 discontinuities, which cross the Sp conversions at the Moho and at the 410. In Fig. 7c ray paths of Sm with a conversion at the receiver-side Moho (Fig. 7b) and of PmSm with an additional P-to-S conversion at the source-side Moho (Fig. 7a) are shown. The PmSm phase is visible in Fig. 7a as a negative precursor of the Moho conversion, which might be mistaken in the real data as a S-to-P conversion at a negative discontinuity below the Moho underneath the station. In receiver functions, however, this phase is reduced due to summation of many events with different source depths and source-side structures. Care should be taken if single seismograms are used because it seems impossible to identify all phases uniquely in these cases.

4 Topography of the discontinuities in the mantle lithosphere

To image the topography and structure of such discontinuities in the study area we plotted in Figs. 8 and 9 several
depth-migrated profiles and in Figs. 10 and 11 several profiles in the time domain. The width of the profiles is two degrees or more of latitude or longitude. The profiles cannot be chosen much narrower because the number of traces available for summation would be too small. Each of the traces in Figs. 10 and 11 was obtained by summation of several hundred traces. The depth domain profiles are chosen along great circles and the time domain data along latitude and longitude. The traces in the time domain profiles (Figs. 10 and 11) are selected by the location of their S-to-P piercing points at 200 km depth. To a first approximation the precursor time may be multiplied by a factor of 10 to obtain the corresponding depth. The IASP91 model is used for depth migration in Figs. 8 and 9. The first and perhaps most important task of the data interpretation is the recognition of patterns in the data. We have chosen variable width of the profiles in the different figures to show the variability of the correlations in dependence on the profile width. In the mantle lithosphere...
below the Moho we observe several clearly correlatable blue signals (velocity decrease downward). We interpreted these signals as being caused by the LAB in the western United States (marked with magenta lines) and the MLD in the central United States (marked with yellow lines) and the cratonic LAB in the eastern United States (marked with green lines). In the following we discuss the topography of each of these negative discontinuities in the mantle lithosphere.

4.1 Structure of the western LAB

North of about 46° N the western LAB extends from the west coast to at least 100° W where it ends (Figs. 8a–d and 9a). In the wider profile of Fig. 9a it even seems to extend to 90° W. It is smoothly dipping from about 100 km depth at the west coast to about 200 km depth at its eastern end. Between 46 and 48° N the depth domain Fig. 8c is not in very good agreement with the time domain Fig. 10 (46–48° N). The reason could be that Fig. 8c has more traces in the north because it follows the great circle. The western LAB is laterally very heterogeneous just north of about 45° N (see Fig. 10; 46–48 and 43–46° N). South of about 45° N the western LAB reaches from about 114 to about 104° W with increasing easterly extension towards the south. There is some indication of a local easterly dip also in this part of the western LAB (see Fig. 8l).

North of about 46° N we have in addition indications of another negative discontinuity closer to the west coast (marked black in Fig. 10). It dips steeply down to about 200 km depth before turning horizontal at this depth. We interpret this structure as the LAB of the Juan de Fuca plate without going into more detail here.
Figure 9. The same data and same correlation marks as in Fig. 8 but summed along broader south–north and west–east profiles. The profile width and orientation is shown in the location maps at the top of the figure. Region (a) (northernmost west–east profile) does not have any stations in Canada. However, due to the shallow incidence angle of the S-receiver functions, the mantle below southern Canada is also sampled by events from the northwest. The westernmost south–north profile (d) shows only the western LAB in the mantle lithosphere. In profile (e) we see at the southern and northern ends the western LAB, and in the central part the deep MLD. In profile (f) we see the cratonic MLD (yellow), the cratonic LAB (green) and in the north the deep western LAB (magenta). In the easternmost south–north profile (g) we see the MLD and weakly the cratonic LAB. See Fig. 1 for explanation of the abbreviations of tectonic units.

4.2 The structure of the MLD

The MLD is marked yellow in the figures and it is visible in most profiles in the central United States near 100 km depth where it is expected. It is poorly seen south of about 37° N, which is probably due to insufficient data (Fig. 8m–p). Between 39 and 45° N (from the northern end of the Colorado Plateau to Yellowstone) the MLD seems to dip from its normal depth near 100 km east of 100° W to about 200 km depth at about 112° W (see Figs. 8e–k and 9b) at the Sevier thrust belt. We note that there is in the area of 40–45° N and 112–106° W no discontinuity near 100 km depth, which could be
an extension of the western LAB or MLD (see Fig. 9b and e). The westerly dip of the MLD is clearer in Fig. 10 (41–43° N). There are similar indications in the profile in Fig. 10 (38–41° N). In Fig. 10 (43–46° N) it may also be visible, although in this profile there is additionally an indication of the east-dipping western LAB. This apparent crossing of discontinuities could be caused by lateral variations within the 3° latitude wide profile. The observation of such a west-dipping structure in the mantle lithosphere in this part of the United States seems to be new to our knowledge. We will return to this question below when Fig. 11 is discussed.

4.3 The structure of the cratonic LAB

The signals we interpreted as cratonic LAB near 200 km depth are marked green in Figs. 8, 9 and 10. These signals are relatively weak in some of the narrow profiles in Fig. 8. In the wider profiles in Fig. 9c, f and g, the cratonic LAB signals are clearer. In the time domain data in Fig. 10 (33–38° N) we see a very strong cratonic LAB separated by a sharp step from the western LAB. In Fig. 10 (38–41° N) the cratonic LAB is also clearly visible. In this 3° latitude wide profile we find also indications of the west-dipping MLD, which suggests strong heterogeneities here. Our observations of the cratonic LAB indicate, together with other S-receiver function observations (Miller and Eaton, 2010, in Canada and Kind et al., 2013, in Scandinavia), that the deep cratonic LAB is visible in converted waves. A study of the complete USArray data in the east is required for a more comprehensive analysis of the cratonic LAB. Close to the east coast the LAB again seems to occur near 100 km depth (Rychert et al., 2007).

4.4 View on the discontinuities with unfiltered S-receiver functions (Fig. 11)

In Fig. 11 the same data are shown as in Fig. 10, except that no filter is applied apart from the deconvolution. The data are much shorter period than in Fig. 10. If a discontinuity is sharp we expect that its signal also becomes sharper if a shorter period filter or no filter is applied. This seems not to be the case for all the upper mantle discontinuities observed in this study such as the western and cratonic LAB and MLD. The discontinuities seem to dissolve in a sequence of sharper discontinuities with a very scattered appearance with sometimes a relatively short correlation length. A kind of scattered lamella structure with stepwise decreasing velocity could cause such signals. This means that these discontinuities are in reality transition zones with stepwise decreasing velocities. A lamella structure of the upper mantle has been observed in long-range controlled source profiles for a long time, for example in Phanerozoic western Europe (see Kind, 1974) and in the cratons of northern Eurasia (see Mechie et al., 1993) and North America (see Thybo and Puchuc, 1997). No phases are marked in Fig.11 in order not to bias the reader. Generally the same phases as in Fig. 10 are also visible in Fig. 11, however, as groups of scattered signals. The signal with the largest amplitudes is the western LAB, especially in Fig. 11 (33–38° N). It spreads out over more than 5 s, which corresponds to at least 50 km. We see in Fig. 11 (48–50° N) the east-dipping western LAB reaching far below the MLD. West-dipping structures are visible in Fig. 11 (43–46° N) and especially clearly in Fig. 11 (41–43° N). The connection between these structures and the cratonic MLD seems obvious. As in the north, indications of an east-dipping structure are also observed in Figs. 8l and 11 (38–41° N) and. The sharp step between the western LAB and cratonic LAB near 105° W is also clear in the unfiltered data in Fig. 11 (33–38° N) as in the longer period filtered data in Fig. 10 (33–38° N).

How do our observations agree with earlier seismic images of the mantle lithosphere in the western and central US? Levander and Miller (2012) also used S-receiver function data from USArray in the same area. They also observe the break in the LAB along the Sevier thrust belt. However they interpret the deep velocity drop west of the Sevier thrust belt as the cratonic LAB of the Laurentia plate. They do not report on the MLD and a west-dipping structure between the Rocky Mountain Front and the Sevier thrust belt. Hopper et al. (2014) observed in the Yellowstone region the same break in the mantle lithosphere, also using S-receiver functions. They observed east of Yellowstone a faint shallow MLD but no west-dipping structure. Hansen et al. (2013) observed the LAB below the Colorado Plateau at 100–150 km depth and east of it at 150–200 km depth. This agrees with our LAB observations along 38° N (Fig. 8l). Foster et al. (2014) also studied the lithosphere in the American Midwest with USArray S-receiver functions. They observed east of about 98° W a strong MLD near 100 km depth and the cratonic LAB at 200–250 km depth. Their data are close to our profile shown in Fig. 8d. We also see here a strong MLD east of about 98° W. However we see, in addition, indications of the east-dipping LAB and the west-dipping MLD within the 2° latitude wide profile in Fig. 8d. Lekic and Fischer (2014) observed below the Colorado Plateau and surroundings scattered negative signals near 100 km depth in S-receiver functions which they interpreted as the LAB beneath the Basin and Range Province and the Rocky Mountain Front, west and east of the Colorado Plateau, respectively. Beneath the stable Colorado Plateau and the Great Plains they interpreted a negative phase in the same depth range as the MLD. We observe below the Colorado Plateau the LAB of the western United States at about 100 km depth (Fig. 9c).

5 Anisotropy of the MLD?

Sodoudi et al. (2013) found in South Africa MLD signals in S-receiver functions showing azimuthal anisotropy. They observed a sign change of the S-to-P converted signal from about 85 km depth with a periodicity of 180° back azimuth.
Figure 10. West–east S-receiver function profiles in the time domain. The latitude of the distribution of piercing points in 200 km depth is marked in each panel. The data are filtered with an 8 s low-pass filter. Correlated phases are marked with the same colors as in Figs. 8 and 9, except that here black represents the LAB of the Juan de Fuca plate.

The MLD in our data does not show any sign change or significant change of the amplitudes as a function of the back azimuth (see LVZ1 in Fig. 5). This means that there are no indications in our data for azimuthal anisotropy as the cause of the MLD signal in the cratonic US. The azimuthal coverage of seismic sources is good for the identification of azimuthal periodicities (see Fig. 2). By far most sources are concentrated in opposing azimuths (Alaska and the Aleutians versus the Andes and the south-western Pacific versus the eastern Mediterranean). The missing anisotropy in MLD S-receiver function data (Fig. 5) seems to be in contradiction to the anisotropy model obtained from a joint inversion of tomography and SKS data (Yuan and Romanowicz, 2010). The reason for this discrepancy is still unknown.

Figure 11. The same profiles as in Fig. 10 but with no filter applied.
could mean that the Archean lithosphere of the craton (Yuan and Romanowicz, 2010) was deformed during the collision with the Farallon plate. The western LAB dips partly far to the east to the mid-continental rift system, where it could be mistaken for the cratonic LAB. The deep cratonic LAB near 200 km depth is weakly observed at the eastern end of the considered area (from about 90°–82° W). Its connection to the previously observed shallow LAB near the Atlantic coast needs further investigation with all eastern USArray data. The cratonic LAB is very strong in the south-western part of the US. Below 200 km depth we have observed a scattered Lehmann discontinuity, which is considered to be the bottom of the asthenosphere. Directly above the 410 km discontinuity mainly in the western United States we observed a strong velocity reduction. Such a velocity reduction directly above the 410 is not observed in the cratonic US, which indicates higher velocity here.

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References

6 Low-velocity zone above the 410 km discontinuity

Tauzin et al. (2010) observed a nearly global velocity reduction at about 350 km depth without any relation to surface tectonics. The cause of this low-velocity layer is thought to be partial melt caused by dehydration (Bercovici and Karato, 2003). We observed a similar discontinuity in the USArray S-receiver functions (see, e.g., Figs. 4 and 9). Jordan and Paulson (2013) suggested that a low-viscosity zone directly above the 410 discontinuity decouples the thick cratonic tectosphere from the flowing mantle. Looking at Fig. 9 it seems that the strongest signals of the phase LVZ410 are observed in the western United States where the LAB occurs at about 100 km depth. The signal is weaker in the cratonic United States (Fig. 9b, c, g). This means that the velocities directly above the 410 are lower in Phanerozoic regions than in cratonic regions of the US. This supports the tomographic results summarized by Jordan and Paulson (2013) of the cratonic tectosphere having higher velocities down to the 410 km discontinuity.

7 Conclusions

We have been able to image with S-receiver functions the major seismic discontinuities in the upper mantle below large regions of the western and central United States. In the upper 200 km we see complex structures of the western LAB and the MLD (Fig. 12). The east-dipping LAB interferes with the west-plunging MLD in a complicated manner. We interpret these structures as being caused by the continental collision of the Farallon plate and the Laurentia plate. The MLD appears to be deformed in this collision in a similar way as could be expected for a shallow non-cratonic LAB. This
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